Numerical methods for time harmonic locally perturbed periodic media

Part 4

Sonia Fliss (Ecole Nationale Supérieure de Techniques Avancées)

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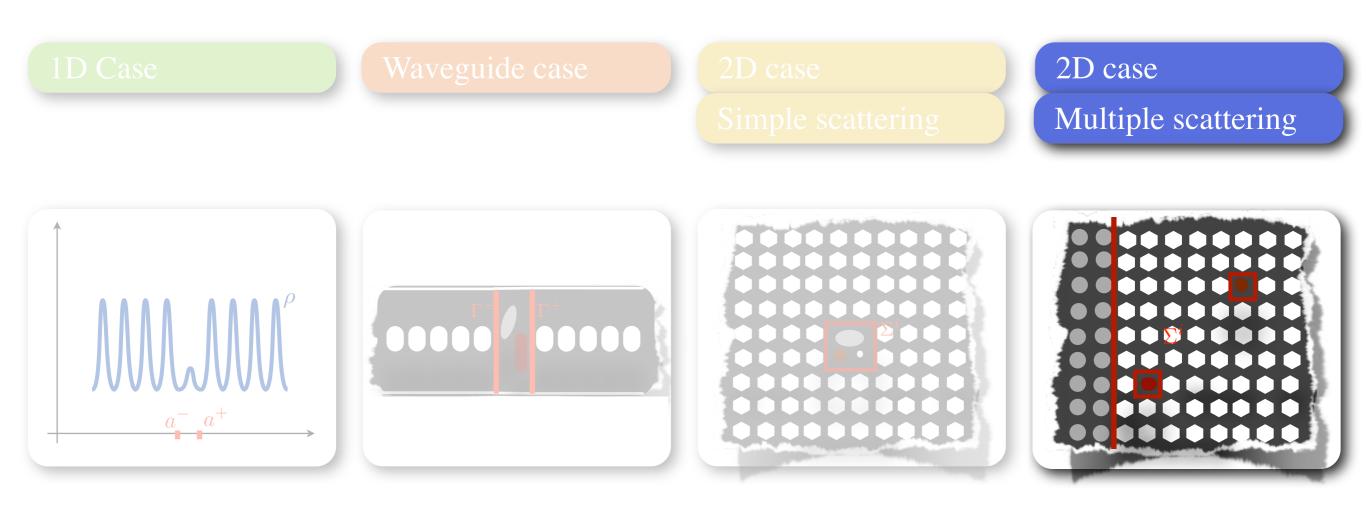
Numerical methods for time harmonic scalar wave equation in locally perturbed periodic media - Part 4

Sonia Fliss

POEMS (UMR 7231 CNRS-INRIA-ENSTA)

Mainly based on joint works with Julien Coatleven, Patrick Joly and Jing-Rebecca Li







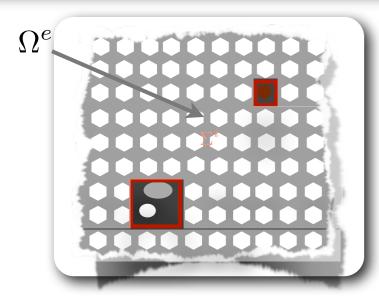
Time harmonic Scalar wave Problem with absorption

$$(\mathcal{P}_{\varepsilon}^{i}) \qquad \begin{vmatrix} -\Delta u_{\varepsilon}^{i} - \rho(\mathbf{x})(\omega^{2} + \imath \varepsilon \omega) u_{\varepsilon}^{i} = f(\mathbf{x}), & \text{in } \Omega^{i} \\ \frac{\partial u_{\varepsilon}^{i}}{\partial \mathbf{n}} + \Lambda_{\varepsilon} u_{\varepsilon}^{i} = 0 & \text{on } \Sigma^{i} = \partial \Omega^{i} \end{vmatrix}$$

$$\Omega^{*} = \Omega_{1}^{*} \cup \Omega_{2}^{*} \quad \Sigma^{*} = \Sigma_{1}^{*} \cup \Sigma_{2}^{*}$$

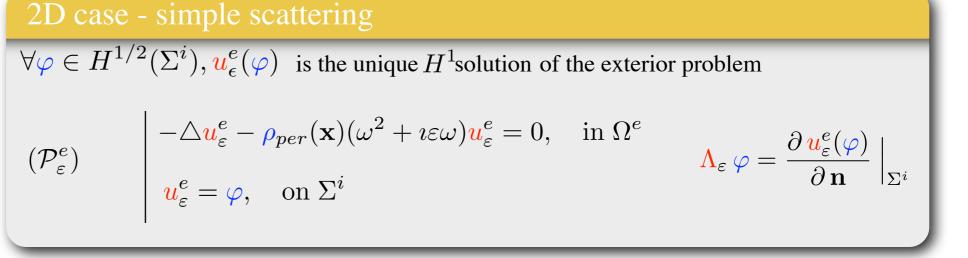
2D case - simple scattering $\forall \varphi \in H^{1/2}(\Sigma^i), u_{\epsilon}^e(\varphi) \text{ is the unique } H^1 \text{ solution}$ of the exterior problem $(\mathcal{P}_{\varepsilon}^e) = -\Delta u_{\varepsilon}^e - \rho_{per}(\mathbf{x})(\omega^2 + \imath \varepsilon \omega) u_{\varepsilon}^e = 0, \text{ in } \Omega^e$ $u_{\varepsilon}^e = \varphi, \text{ on } \Sigma^i$ $\Lambda_{\varepsilon} \varphi = \frac{\partial u_{\varepsilon}^e(\varphi)}{\partial \mathbf{n}} \Big|_{\Sigma^i}$





 Ω_1^e

 $\Omega_{2^*}^e$



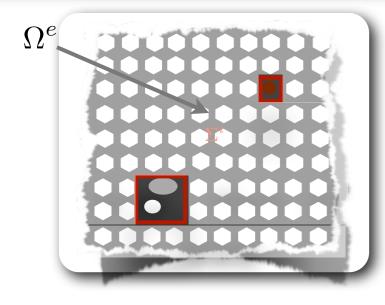
 $\forall \varphi_1 \in H^{1/2}(\Sigma_1^i), u_1^e(\varphi_1) \text{ is the unique } H^1 \text{ solution of the exterior problem posed in } \Omega_1^e \\ \Lambda_1 \text{ corresponding DtN operator}$

Theorem

If $\Sigma_1^i \cap \Sigma_2^i = \emptyset$, $\forall \varphi \in H^{1/2}(\Sigma^i)$, $\exists ! (\varphi_1, \varphi_2) \in H^{1/2}(\Sigma_1^i) \times H^{1/2}(\Sigma_2^i)$ $\frac{u^e(\varphi) = u_1^e(\varphi_1) \Big|_{\Omega^e} + u_2^e(\varphi_2) \Big|_{\Omega^e}$

 $\forall \varphi_2 \in H^{1/2}(\Sigma_2^i), u_2^e(\varphi_2)$ is the unique H^1 solution of the exterior problem posed in Ω_2^e Λ_2 corresponding DtN operator





 Ω_1^e

 Ω_2^e



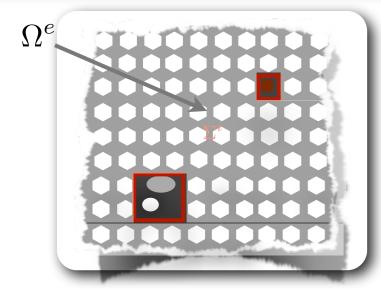
 $\forall \varphi \in H^{1/2}(\Sigma^i), {\pmb u}^e_\epsilon(\varphi) \ \text{ is the unique } H^1 \text{ solution of the exterior problem }$

$$(\mathcal{P}^{e}_{\varepsilon}) \qquad \begin{vmatrix} -\Delta \boldsymbol{u}^{e}_{\varepsilon} - \rho_{per}(\mathbf{x})(\omega^{2} + \imath\varepsilon\omega)\boldsymbol{u}^{e}_{\varepsilon} = 0, & \text{in } \Omega^{e} \\ \boldsymbol{u}^{e}_{\varepsilon} = \varphi, & \text{on } \Sigma^{i} \end{vmatrix} \qquad \qquad \Lambda_{\varepsilon} \varphi = \frac{\partial \boldsymbol{u}^{e}_{\varepsilon}(\varphi)}{\partial \mathbf{n}} \Big|_{\Sigma^{i}}$$

Theorem

If
$$\Sigma_1^i \cap \Sigma_2^i = \emptyset$$
, $\forall \varphi \in H^{1/2}(\Sigma^i), \exists !(\varphi_1, \varphi_2) \in H^{1/2}(\Sigma_1^i) \times H^{1/2}(\Sigma_2^i)$
$$\frac{\mathbf{u}^e(\varphi) = \mathbf{u}_1^e(\varphi_1) \Big|_{\Omega^e} + \mathbf{u}_2^e(\varphi_2) \Big|_{\Omega^e}$$





 Ω_1^e

 Ω_2^e



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Theorem

 $21 \varphi_1$

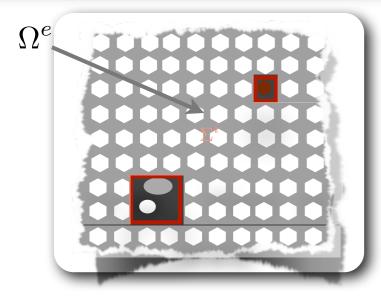
If
$$\Sigma_1^i \cap \Sigma_2^i = \emptyset$$
, $\forall \varphi \in H^{1/2}(\Sigma^i), \exists ! (\varphi_1, \varphi_2) \in H^{1/2}(\Sigma_1^i) \times H^{1/2}(\Sigma_2^i)$
$$\frac{u^e(\varphi) = u_1^e(\varphi_1) \Big|_{\Omega^e} + u_2^e(\varphi_2) \Big|_{\Omega^e}$$

$$\begin{split} \left| \begin{array}{c} \Lambda_{\varepsilon} \varphi \right|_{\Sigma_{1}^{i}} &= \Lambda_{1} \varphi_{1} + \frac{\partial}{\partial \mathbf{n}} u_{2}^{e}(\varphi_{2}) \Big|_{\Sigma_{1}^{i}} \\ & & & \\ \Lambda_{12} \varphi_{2} \end{array} \right| \quad \begin{pmatrix} \Lambda_{\varepsilon} \\ \Lambda_{12} \varphi_{2} \\ & \end{pmatrix} \\ \left| \begin{array}{c} \Lambda_{\varepsilon} \varphi \\ \Lambda_{\varepsilon} \varphi \\ & \\ \Sigma_{2}^{i} \end{array} \right|_{\Sigma_{2}^{i}} = \frac{\partial}{\partial \mathbf{n}} u_{1}^{e}(\varphi_{1}) \Big|_{\Sigma_{2}^{i}} + \Lambda_{2} \varphi_{2} \\ & & \\ & & \\ & & \\ & & \\ & & \\ \end{array} \right| \quad \text{wh} \quad \end{split}$$

$$\begin{pmatrix} \Lambda_{\varepsilon} \varphi \Big|_{\Sigma_{1}^{i}}, \Lambda_{\varepsilon} \varphi \Big|_{\Sigma_{2}^{i}} \end{pmatrix} = \Lambda \begin{pmatrix} \varphi_{1}, \varphi_{2} \end{pmatrix}$$
where $\Lambda = \begin{bmatrix} \Lambda_{1} & \Lambda_{12} \\ \Lambda_{21} & \Lambda_{2} \end{bmatrix}$



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 Ω_1^e



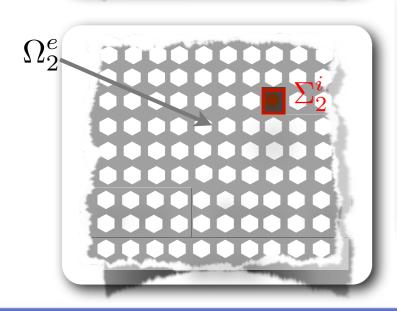
 $\forall \varphi \in H^{1/2}(\Sigma^i), \mathbf{u}^e_\epsilon(\varphi)$ is the unique H^1 solution of the exterior problem

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Theorem

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$$\frac{u^e(\varphi) = u_1^e(\varphi_1) \Big|_{\Omega^e} + u_2^e(\varphi_2) \Big|_{\Omega^e}$$

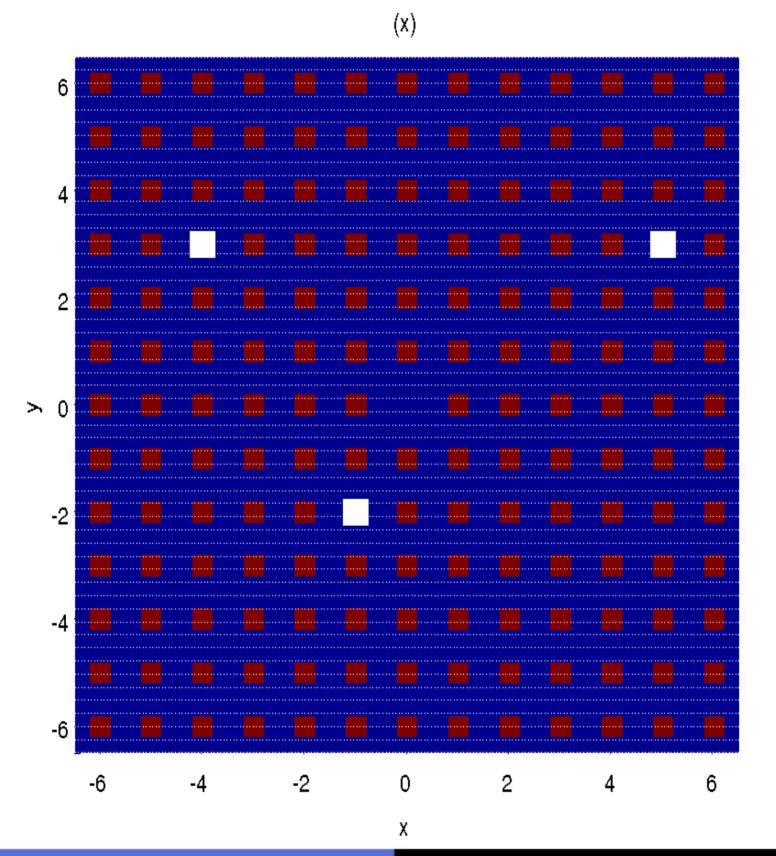
Theorem



where
$$\Theta = \begin{bmatrix} \mathbb{I} & \Theta_{12} \\ \Theta_{21} & \mathbb{I} \end{bmatrix}$$
 and $\Lambda = \begin{bmatrix} \Lambda_1 & \Lambda_{12} \\ \Lambda_{21} & \Lambda_2 \end{bmatrix}$
For homogenous media : Balabane & Tirel (1997), Grote & Kirsch (2004), Ben Hassen et al. (2007)

Numerical results



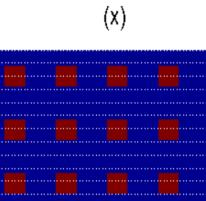


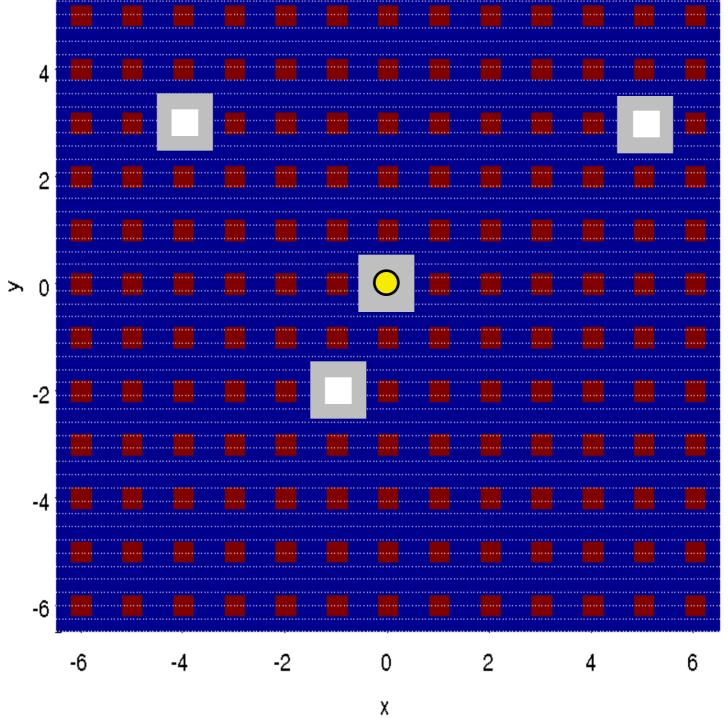
Sonia Fliss Exact boundary conditions, Wave propagation and Periodic media

Numerical results

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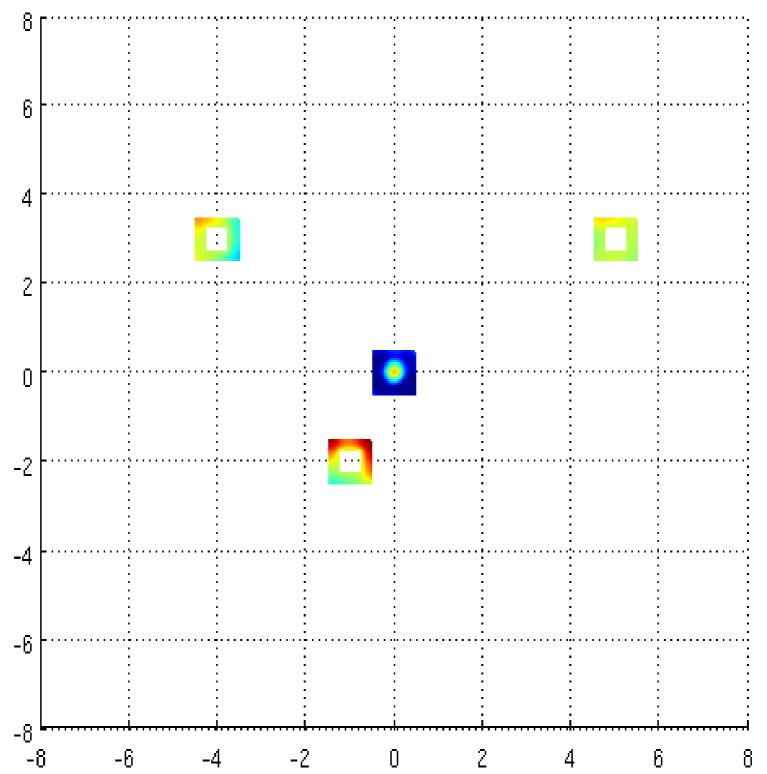




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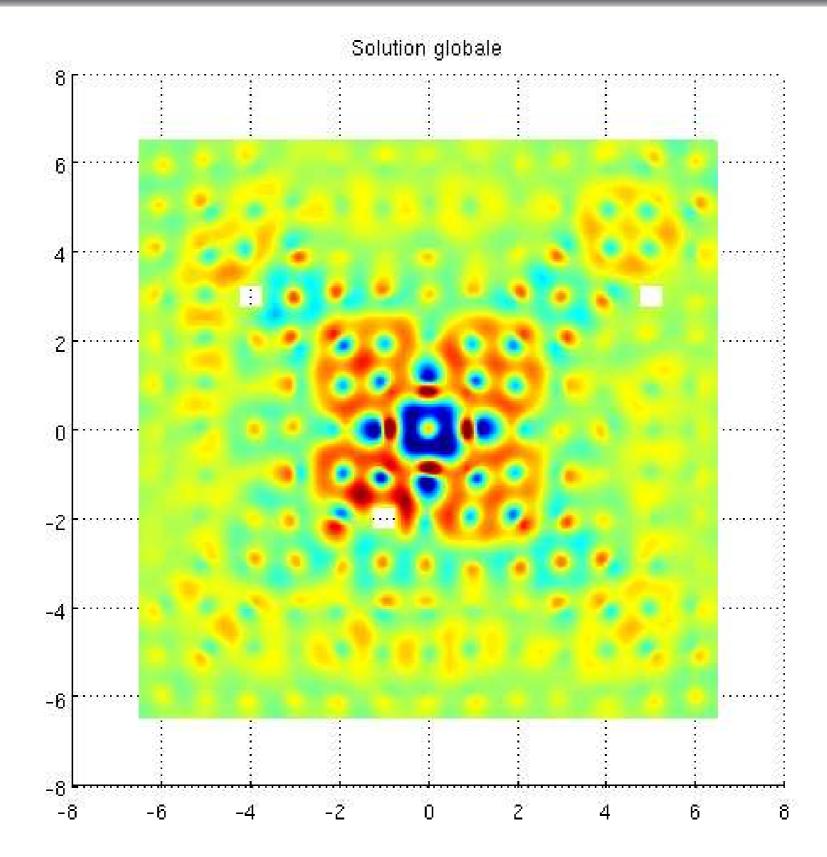


Solution interieure



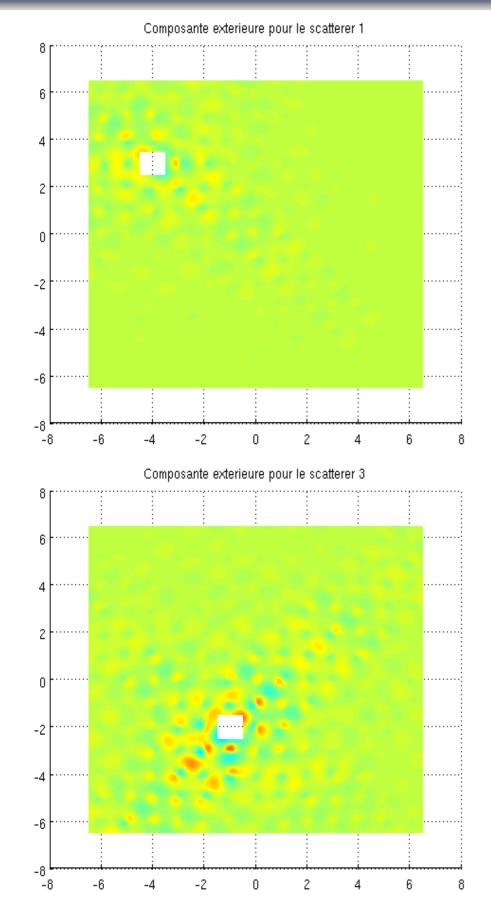
Numerical results



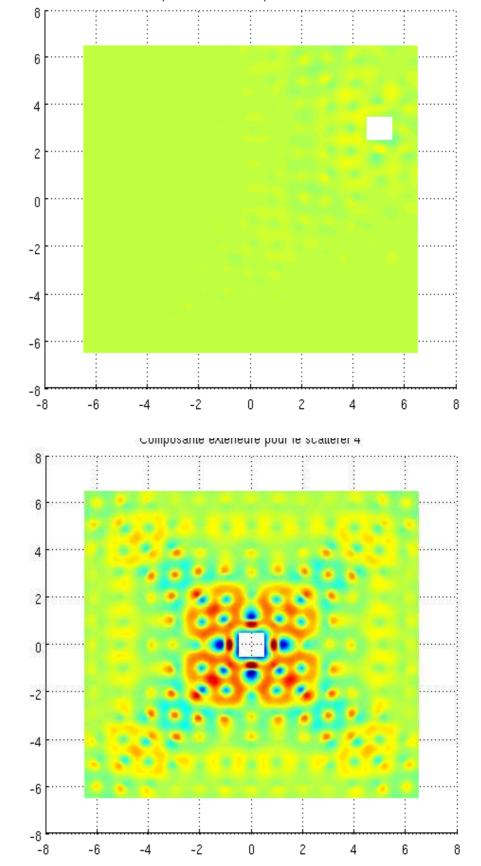


Numerical results



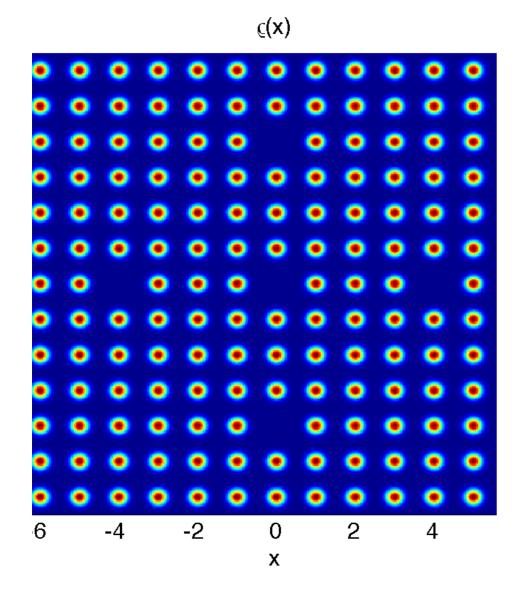


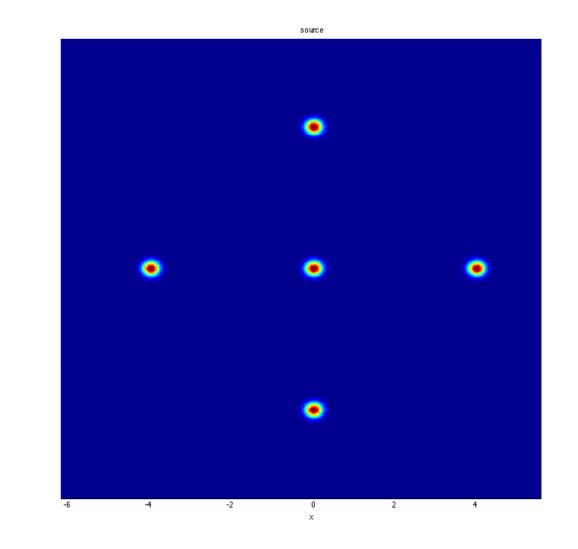
Composante exterieure pour le scatterer 2



Poems

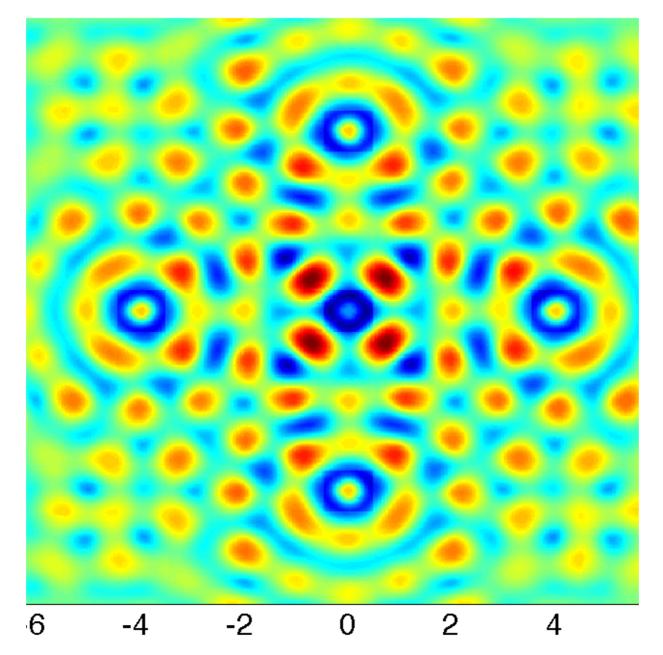
The 2D plane problem - Multiple scattering



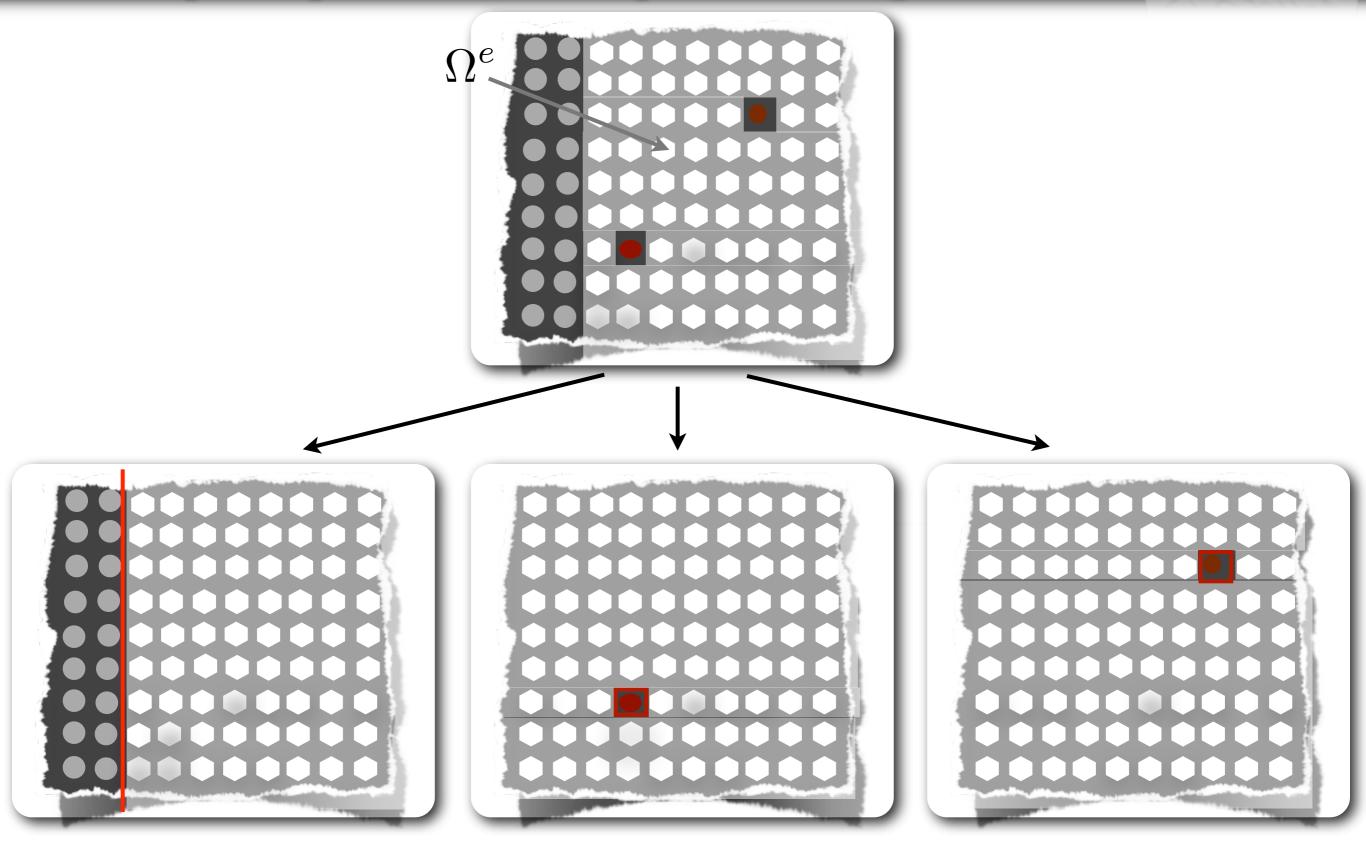




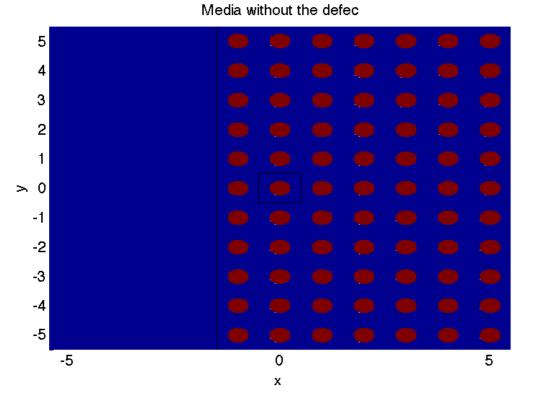
Solution globale



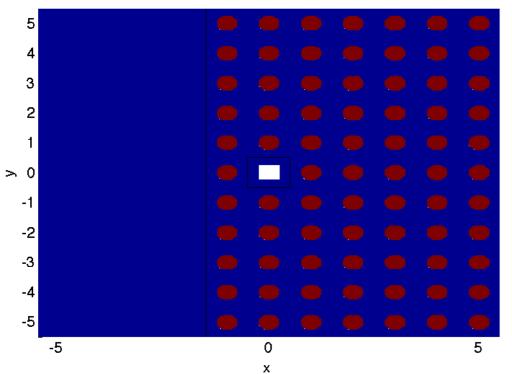




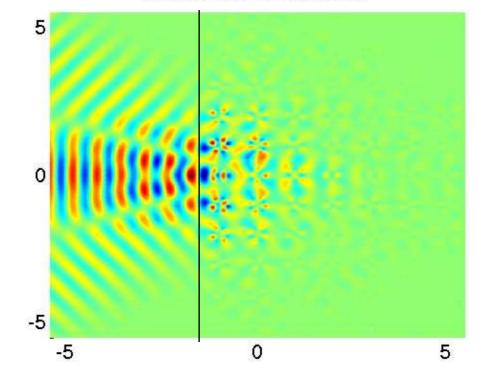




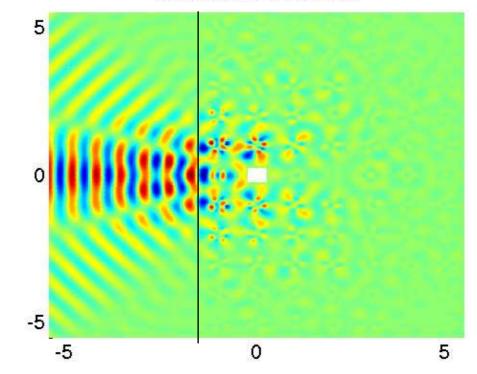
Media with the defec



Solution without the defec

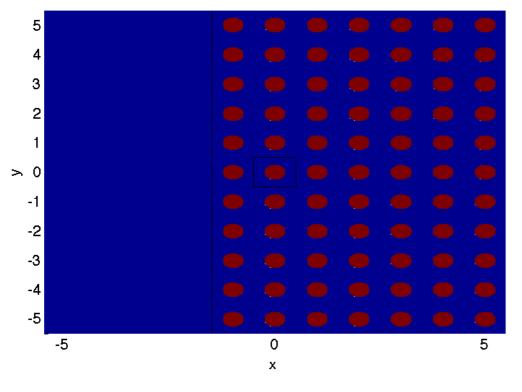


Solution with the defect

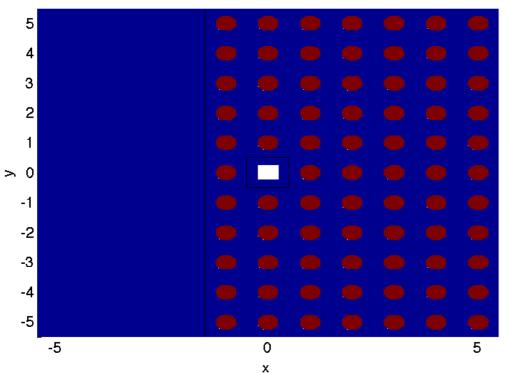




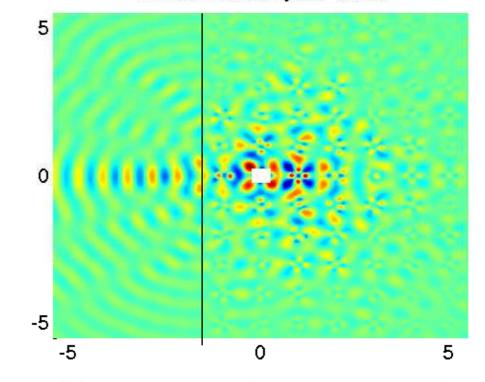
Media without the defec



Media with the defec



Wave diffracted by the defect





Numerical methods for time harmonic scalar wave equation in locally perturbed periodic media - Part 4

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Numerical methods for time harmonic scalar wave equation in locally perturbed periodic media - Conclusions

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Absorbing media

The theory is complete and the numerical method works well

Non absorbing media

Limiting absorption principle for the definition of the DtN operators

Under some conditions on the periodic media, the problem is Fredholm

The problem is well posed except for a countable set of frequencies?



Absorbing media

The theory is complete and the numerical method works well

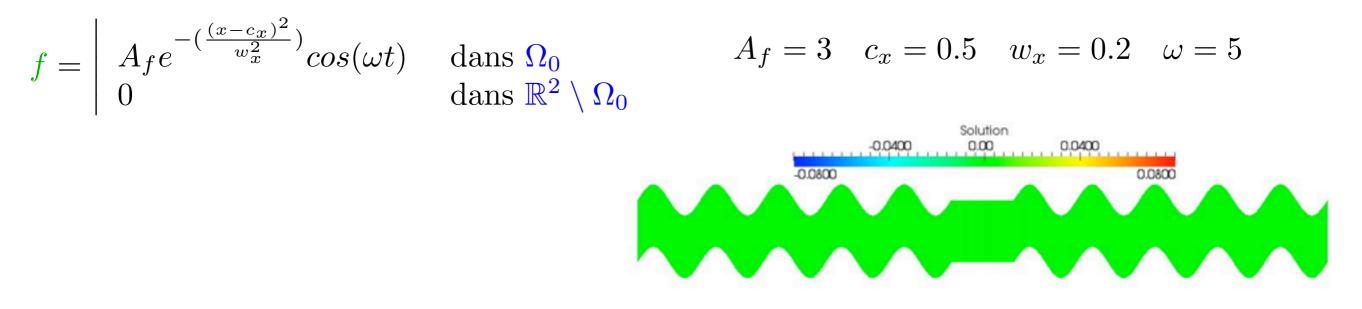
The numerical analysis has to be done

Non absorbing media

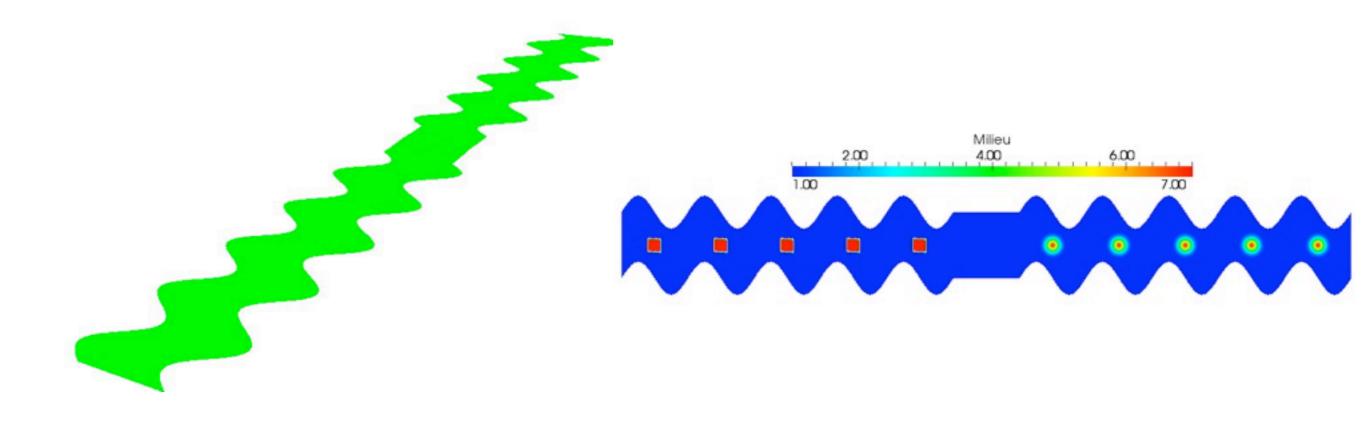
A numerical limiting absorption method is done for the other cases

The corresponding theory still raises challenging open questions

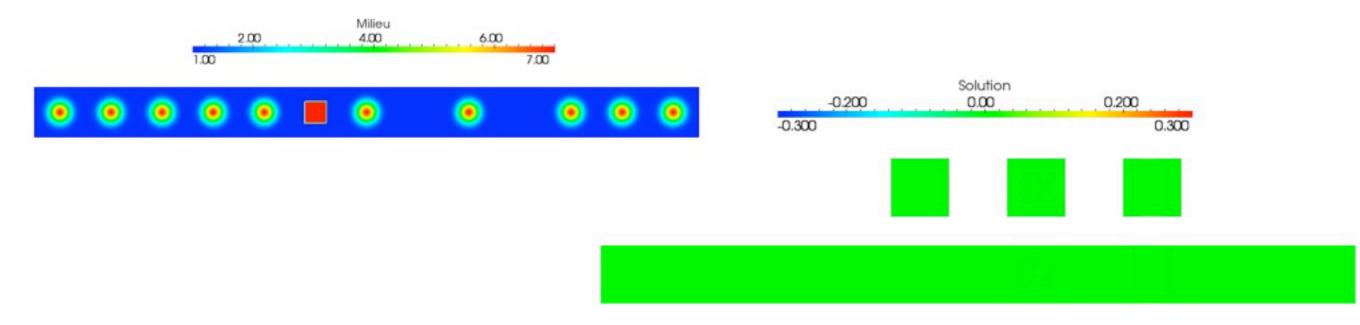




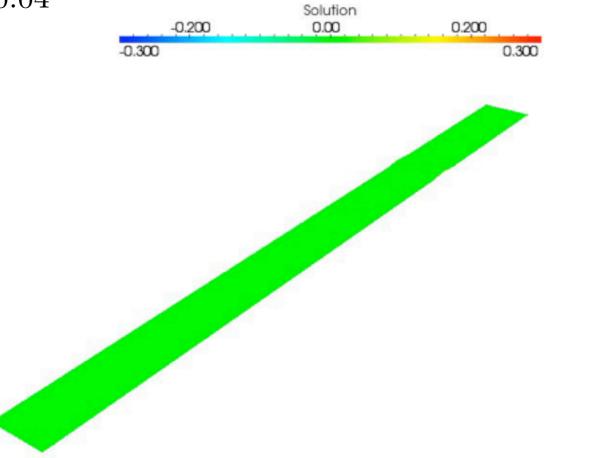
 $r = 2, \ \theta = 1/4, \ \Delta t = 0.04$



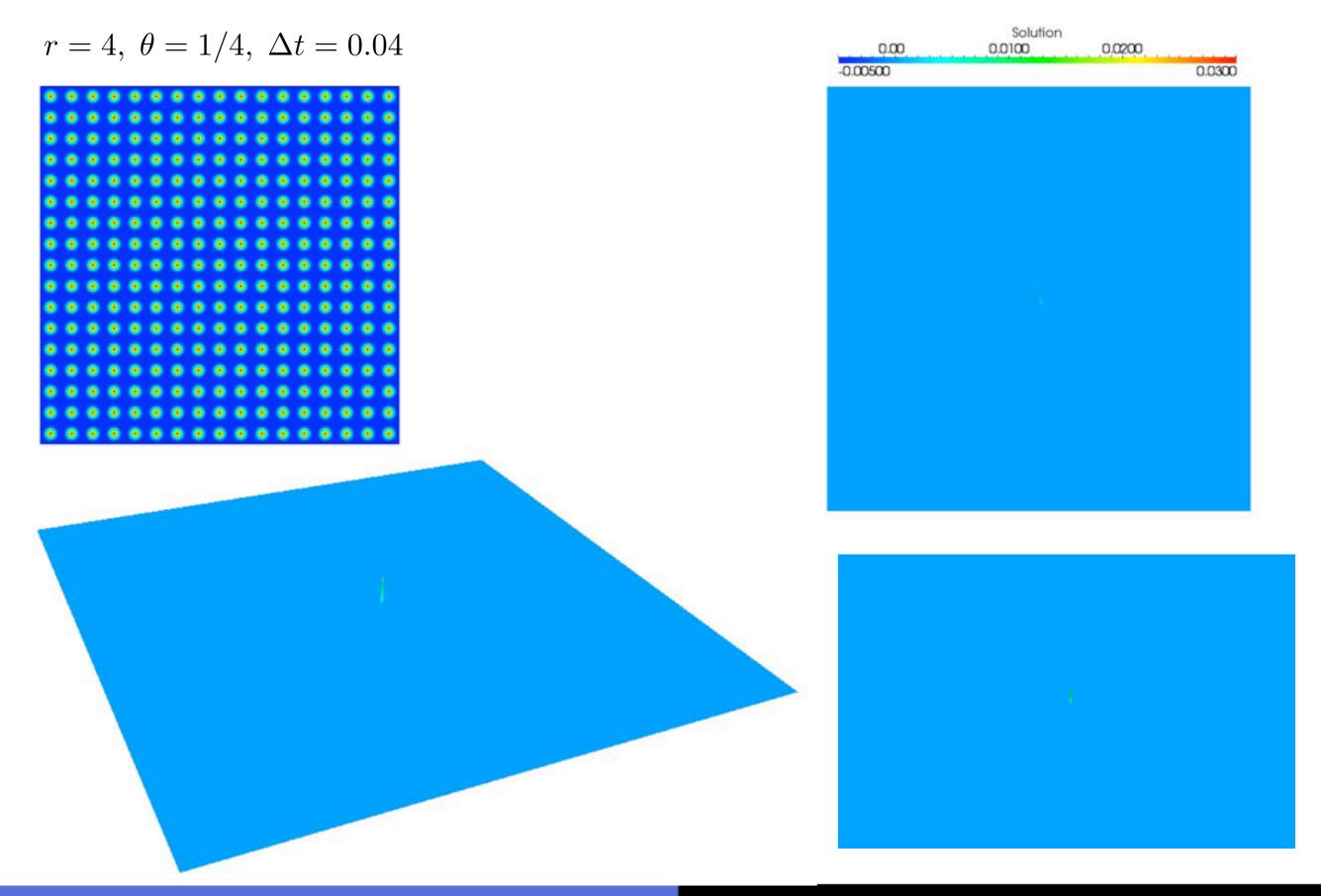




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r = 3, \ \theta = 1/4, \ \Delta t = 0.04
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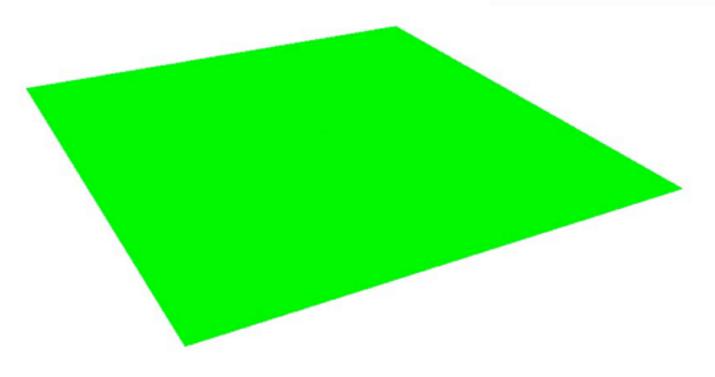








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Thank you for your attention